

Methods for Computational Fluid Dynamics Aiming Quantum Computing

Thiago F. P. O. Chahin¹, Leonardo R. Monteiro²
LACIA/UDESC, Joinville, SC

This work investigates linearization techniques for the Navier-Stokes equations (NSE) using the Taylor-Green Vortex (TGV) benchmark, with applications to Quantum Computing (QC). Pre-processing NSE for QC will open avenues for faster fluid simulations in the future [4].

Quantum algorithms require linear formulations, as quantum operations follow the superposition principle through unitary transformations [2]. This fundament is shown in Dirac notation:

$$U(a|\psi_1\rangle + b|\psi_2\rangle) = aU|\psi_1\rangle + bU|\psi_2\rangle, \quad (1)$$

where U is a unitary operator ($U^\dagger U = I$), $|\psi_i\rangle$ represent quantum state vectors, and $a, b \in \mathbb{C}$ are complex probability amplitudes. The nonlinear convective term $(\mathbf{u} \cdot \nabla)\mathbf{u}$ in NSE violates this linearity requirement, necessitating specialized approximation techniques. The TGV problem provides an ideal test case for its exact analytical solution and periodic boundary conditions [1, 7].

The TGV test case is governed by the incompressible NSE [7]:

$$\frac{\partial \mathbf{u}}{\partial t} + (\mathbf{u} \cdot \nabla)\mathbf{u} = -\nabla p + \nu \nabla^2 \mathbf{u}, \quad \nabla \cdot \mathbf{u} = 0, \quad (2)$$

where \mathbf{u} is the velocity field, p is the pressure, ν is the kinematic viscosity, and t is time.

The analytical solution for TGV is [7]:

$$\mathbf{u}_{\text{analytical}} = \begin{pmatrix} A \cos(ax) \sin(by) \sin(cz) e^{-\nu(a^2+b^2+c^2)t} \\ B \sin(ax) \cos(by) \sin(cz) e^{-\nu(a^2+b^2+c^2)t} \\ C \sin(ax) \sin(by) \cos(cz) e^{-\nu(a^2+b^2+c^2)t} \end{pmatrix}, \quad (3)$$

where A, B, C, a, b , and c are constants, and x, y , and z are Cartesian coordinates.

The studied methods are:

1. **Local Temporal Linearization:** Approximates the nonlinear term using a Taylor series expansion around equilibrium \mathbf{u}_0 [6]:

$$(\mathbf{u} \cdot \nabla)\mathbf{u} \approx (\mathbf{u}_0 \cdot \nabla)\mathbf{u}_0 + \alpha [(\mathbf{u}_0 \cdot \nabla)\delta\mathbf{u} + (\delta\mathbf{u} \cdot \nabla)\mathbf{u}_0], \quad (4)$$

where $\delta\mathbf{u} = \mathbf{u} - \mathbf{u}_0$ is a small perturbation, and α is a coefficient that calibrates the velocity field at each time step.

2. **SVD Matricial Tensorial Linearization:** Uses Singular Value Decomposition (SVD) to approximate velocity fields in low-rank format, reducing complexity while preserving flow features. The 3D fields are reshaped into 2D matrices, decomposed via SVD, and rebuilt using dominant singular components:

$$\mathbf{u}_{\text{approx}} = \mathbf{U}_{:,1:r} \cdot \mathbf{S}_{1:r,1:r} \cdot \mathbf{V}_{1:r,:}^T, \quad (5)$$

¹thiago.chahin@edu.udesc.br

²leonardo.monteiro@udesc.br

where \mathbf{U} , \mathbf{S} , and \mathbf{V} are singular vectors and values, and r is the rank. The linearized fields are updated as:

$$\mathbf{u}_{\text{linearized}} = \mathbf{u}_{\text{approx}} - \alpha \mathbf{u}_{\text{approx}} \Delta t, \quad (6)$$

where α is a linearization coefficient and Δt is the time step. This approach aligns with low-rank solvers for Navier-Stokes equations [3].

3. Logarithmic Linearization: This method linearizes nonlinear velocity terms by applying a logarithmic transformation to $|\mathbf{u}|$, ensuring positivity with a small constant $\epsilon = 10^{-10}$. Linearization is performed in log-space and mapped back via the exponential function, preserving velocity direction [5]:

$$\mathbf{u}_{\log} = \log(|\mathbf{u}| + \epsilon), \quad \mathbf{u}_{\text{linearized}} = \exp(\mathbf{u}_{\log} - \alpha \mathbf{u}_{\log} \Delta t) \cdot \frac{\mathbf{u}}{|\mathbf{u}|}, \quad (7)$$

where α is a linearization coefficient and Δt the time step.

Numerical experiments used a 64^3 grid, 2π domain, $\nu = 0.01 \text{ m}^2/\text{s}$, $\Delta t = 0.001 \text{ s}$, third-order Runge-Kutta time integration [1], and finite-differences for spatial derivatives.

The methods yielded comparable accuracy to standard NSE:

- **Navier-Stokes:** MSE = 0.049062, Absolute Error = 0.161511
- **Local Temporal:** MSE = 0.049062, Absolute Error = 0.161511
- **SVD:** MSE = 0.049001, Absolute Error = 0.161220
- **Logarithmic:** MSE = 0.049049, Absolute Error = 0.161493

The SVD method showed superior performance, demonstrating potential for quantum computing applications where linear formulations are essential.

References

- [1] M. E. Brachet, D. I. Meiron, S. A. Orszag, B. G. Nickel, R. H. Morf, and U. Frisch. “Small-scale structure of the Taylor-Green vortex”. In: *Journal of Fluid Mechanics* 130 (1983), pp. 411–452. DOI: 10.1017/S0022112083001159.
- [2] A. W. Harrow. “Quantum Algorithms for Systems of Linear Equations”. In: *arXiv preprint* (2015). <https://arxiv.org/abs/1501.0008>.
- [3] K. Lee, H. C. Elman, and B. Sousedík. “A Low-Rank Solver for the Navier-Stokes Equations with Uncertain Viscosity”. In: *SIAM Journal on Scientific Computing* 41.5 (2019), A1277–A1304. DOI: 10.1137/18M1186801.
- [4] X. Li, X. Yin, N. Wiebe, J. Chun, G. K. Schenter, M. S. Cheung, and J. Millmenstadt. “Potential quantum advantage for simulation of fluid dynamics”. In: *Physical Review Research* 7.1 (2025), p. 013036. DOI: 10.1103/PhysRevResearch.7.013036. URL: <https://journals.aps.org/prresearch/abstract/10.1103/PhysRevResearch.7.013036>.
- [5] M. Ltifi. “Strong solution of modified 3D-Navier-Stokes equations”. In: *arXiv preprint* (2021). <https://arxiv.org/abs/2111.00859>.
- [6] T. W. H. Sheu and R. K. Lin. “Newton linearization of the incompressible Navier-Stokes equations”. In: *International Journal for Numerical Methods in Fluids* 44.3 (2004), pp. 297–312. DOI: 10.1002/flid.639.
- [7] G. I. Taylor and A. E. Green. “Mechanism of the production of small eddies from large ones”. In: *Proceedings of the Royal Society of London. Series A, Mathematical and Physical Sciences* 158.895 (1937), pp. 499–521. DOI: 10.1098/rspa.1937.0035.